

# Degeneracy and first order phase transitions

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Des Johnston (bad guy)  
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# Plan of talk

Mea Culpa

A family of  $3D$  Ising models motivated by random surface simulations

A problem with the 1st order FSS for a plaquette  $3D$  Ising model in this family

A solution

# Mea Culpa I

Wolfhard has been interested in first order transitions for a very long time

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## FIRST ORDER IN 2D DISCLINATION MELTING

W. JANKE and H. KLEINERT

*Institut für Theorie der Elementarteilchen, Freie Universität Berlin,  
Arnimallee 14, 1000 Berlin 33, West-Germany*

Received 24 April 1984

We present a Monte Carlo study of a 2D gas of crystal disclinations and find a first-order melting transition.

Prejudice: first order transitions poor cousins of second order. Fewer different numbers to measure....and harder to do so.

# Mea Culpa II

Even sharing an office with Wolfhard in 1989 wasn't enough to change my opinion

Though there were some distractions.....



Only enlightened by recent work

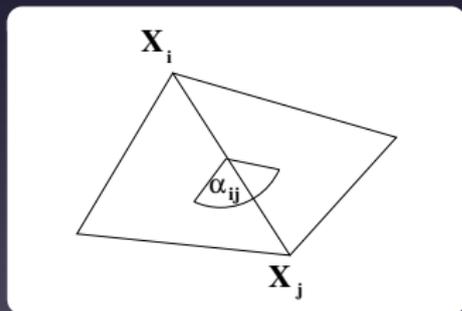
# The Hamiltonian(s) of Interest

# The Gonihedric action

Savvidy “Gonihedric” Hamiltonian

$$H = \sum_{ij} |X^\mu(i) - X^\mu(j)| \theta_{ij}$$
$$\theta_{ij} = ||\pi - \alpha_{ij}||$$

Gonia: angle  
Hedra: face

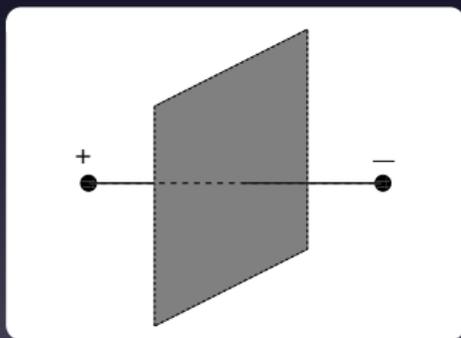


# Spins Cluster Boundaries as Surface Models

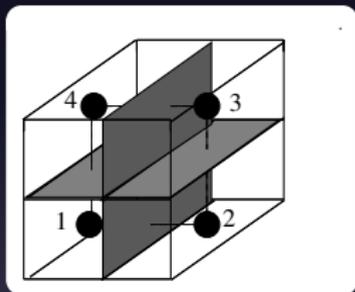
Spin cluster boundaries  $\leftrightarrow$  surfaces

Edge spins:  $U_{ij} = -1$

Vertex spins:  $\sigma_i \sigma_j = -1$



# Counting configurations with spins (areas and intersections)



# Ising/Surface correspondence

Allow energy from areas, edges and intersections (A. Cappi, P Colangelo, G. Gonella and A. Maritan)

$$H = \sum (\beta_A n_A + \beta_C n_C + \beta_I n_I)$$

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$$\beta_A = 2J_1 + 8J_2, \quad \beta_C = 2J_3 - 2J_2, \quad \beta_I = -4J_2 - 4J_3$$

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$$H = J_1 \sum_{\langle ij \rangle} \sigma_i \sigma_j + J_2 \sum_{\langle\langle ij \rangle\rangle} \sigma_i \sigma_j + J_3 \sum_{[i,j,k,l]} \sigma_i \sigma_j \sigma_k \sigma_l$$

# Gonihedric $\rightarrow$ Tune Out Area Term

One parameter family of “Gonihedric” Ising models  
(Savvidy, Wegner)

$$\mathcal{H}^\kappa = -2\kappa \sum_{\langle i,j \rangle} \sigma_i \sigma_j + \frac{\kappa}{2} \sum_{\langle\langle i,j \rangle\rangle} \sigma_i \sigma_j - \frac{1-\kappa}{2} \sum_{\square} \sigma_i \sigma_j \sigma_k \sigma_l$$

$\kappa = 0$  pure plaquette

$$\mathcal{H} = - \sum_{\square} \sigma_i \sigma_j \sigma_k \sigma_l$$

# The Dual ( $\kappa = 0$ )

An anisotropically coupled Ashkin-Teller model

$$\mathcal{H}_{dual} = - \sum_{\langle ij \rangle_x} \sigma_i \sigma_j - \sum_{\langle ij \rangle_y} \tau_i \tau_j - \sum_{\langle ij \rangle_z} \sigma_i \sigma_j \tau_i \tau_j,$$

Standard duality relation

$$\tanh \beta = e^{-2\beta^*}$$

# Known: First Order Transition

Strong first order transition

$$\mathcal{H} = - \sum_{\square} \sigma_i \sigma_j \sigma_k \sigma_l$$

Only inaccurate (Metropolis, yours truly, bad ...)  
determination of transition point

Same for dual model - (Metropolis, yours truly, bad ...)  
determination of transition point

# “Standard” First Order Transitions

# The $q$ -state Potts model

Hamiltonian

$$\mathcal{H}_q = - \sum_{\langle ij \rangle} \delta_{\sigma_i, \sigma_j}$$

Evaluate the partition function, derivatives give observables

$$Z(\beta) = \sum_{\{\sigma\}} \exp(-\beta \mathcal{H}_q)$$

# 1st Order FSS: Heuristic two-phase model

A fraction  $W_o$  in  $q$  ordered phase(s), energy  $e_o$

A fraction  $W_d = 1 - W_o$  in disordered phase, energy  $e_d$

Ignore transits

# 1st Order FSS: Energy moments

Energy moments become

$$\langle e^n \rangle = W_o e_o^n + (1 - W_o) e_d^n$$

And the specific heat then reads:

$$C_V(\beta, L) = L^d \beta^2 \left( \langle e^2 \rangle - \langle e \rangle^2 \right) = L^d \beta^2 W_o (1 - W_o) \Delta e^2$$

Max of  $C_V^{\max} = L^d (\beta^\infty \Delta e / 2)^2$  at  $W_o = W_d = 0.5$

Volume scaling

# 1st Order FSS: Specific Heat peak shift

Probability of being in any of the states

$$W_o \sim q \exp(-\beta L^d f_o), \quad W_d \sim \exp(-\beta L^d f_d)$$

Take logs, expand around  $\beta^\infty$

$$\begin{aligned} \ln(W_o/W_d) &= \ln q + \beta L^d (f_d - f_o) \\ &= \ln q + L^d \Delta e (\beta - \beta^\infty) \end{aligned}$$

Solve for specific heat peak  $W_o = W_d, \ln(W_o/W_d) = 0$

$$\beta^{C_V^{\max}}(L) = \beta^\infty - \frac{\ln q}{L^d \Delta e} + \dots$$

# 1st Order FSS: summary

Peaks grow as  $L^d$

Critical points shift as  $1/L^d$

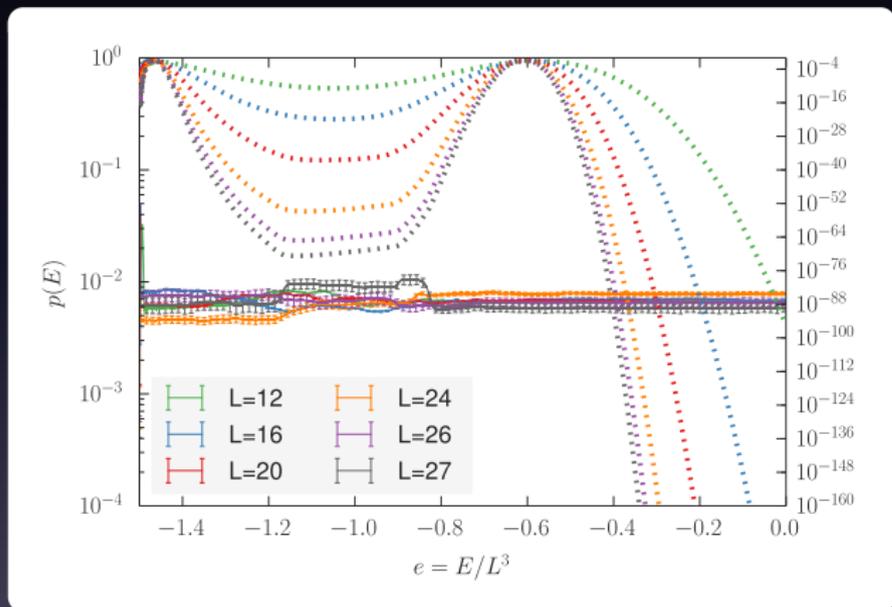
Except ...

Fixed boundaries ( $1/L$  leading term)

$$Z(\beta) = \left[ e^{-\beta(L^d f_d + L^{d-1} \tilde{f}_o)} + q e^{-\beta(L^d f_o + L^{d-1} \tilde{f}_d)} \right] [1 + \dots]$$

# A Problem

# Plaquette Model: Careful Multicanonical Simulations



Multicanonical histograms

# Scaling of the Plaquette Model and Dual

Determine critical point(s)  $L = 8 \dots 27$ , periodic bc,  $1/L^3$  fits - a nice exercise for a PhD student (Marco)

Original model:

$$\beta^\infty = 0.549994(30)$$

Dual model:

$$\beta_{dual}^\infty = 1.31029(19)$$

$$\beta^\infty = 0.55317(11)$$

**Estimates are about 30 error bars apart**

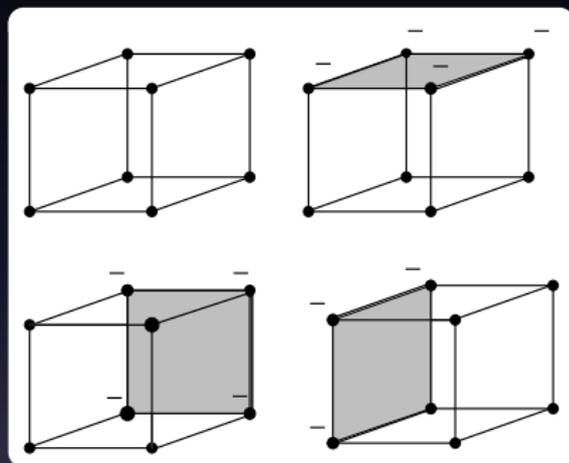
# Potential Solutions

Blame the student (yours truly, bad....)

Incorrect, try again (good guys)

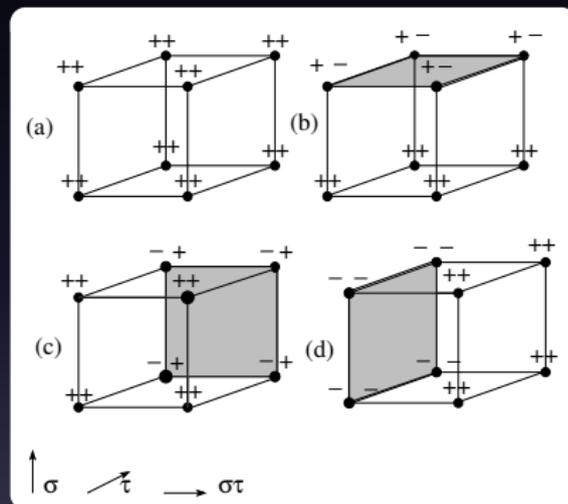
What is special about plaquette model?

# Groundstates: Plaquette



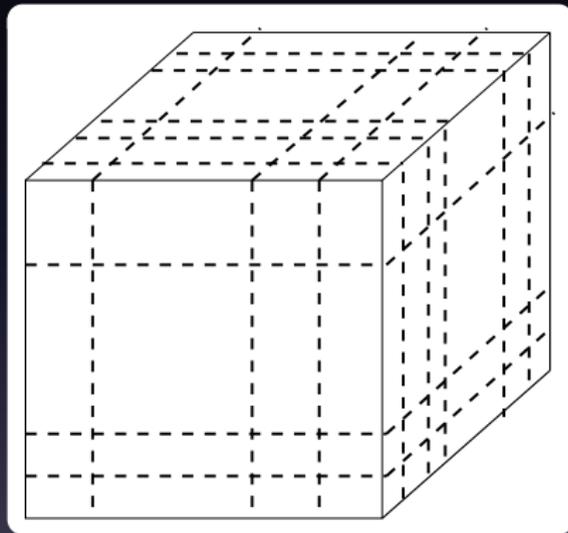
Persists into low temperature phase: degeneracy  $2^{3L}$

# Groundstates: Dual



Dual degeneracy

# Ground state



# 1st Order FSS with Exponential Degeneracy

Normally  $q$  is constant

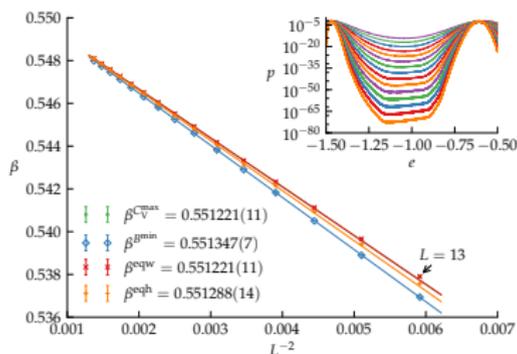
Suppose instead  $q \propto e^L$

$$\beta^{C_V^{\max}}(L) = \beta^\infty - \frac{\ln q}{L^d \Delta e} + \dots$$

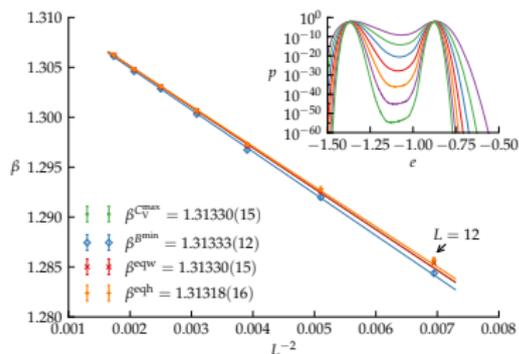
becomes

$$\beta^{C_V^{\max}}(L) = \beta^\infty - \frac{1}{L^{d-1} \Delta e} + \dots$$

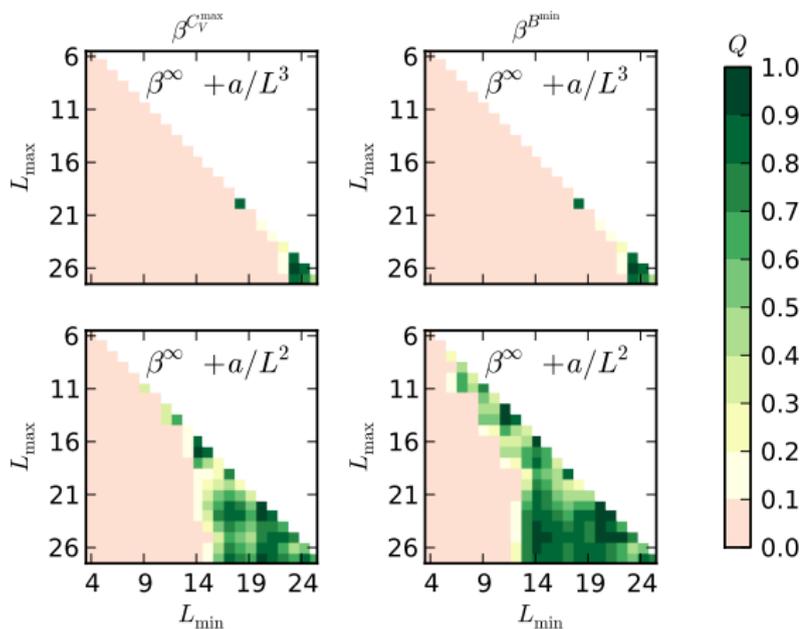
## Plaquette Hamiltonian fits



## Dual Hamiltonian fits



# Quality of fits



Forcing a fit to  $1/L^3$  gives much poorer quality

# Conclusions

Standard 1st order FSS:  $1/L^3$  corrections in 3D

Fixed BC:  $1/L$  (surface tension)

Exponential degeneracy:  $1/L^2$  in 3D

# References

G.K. Savvidy and F.J. Wegner, Nucl. Phys. B **413**, 605 (1994).

M. Mueller, W. Janke and D. A. Johnston,  
Phys. Rev. Lett. **112** (2014) 200601.

M. Mueller, D. A. Johnston and W. Janke, Nucl. Phys. B **888** (2014) 214; Nucl. Phys. B **894** (2015) 1.

Marco's poster

