

Macroscopic Degeneracy and Order in a 3D Plaquette Ising model

Marco Mueller, Wolfhard Janke, Des Johnston
SPLDS, Pont-a-Mousson, May 2015

Plan of talk

A 3D plaquette Ising model with a highly degenerate low-T phase

Consequences for FSS at first order transition

Consequences for order parameter

A 3D Plaquette Ising action

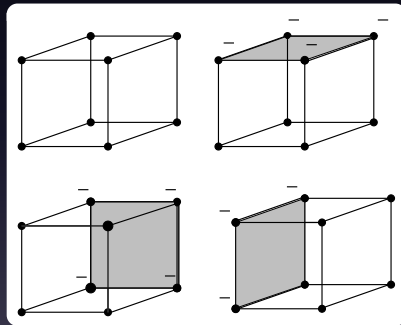
3D cubic lattice, spins on *vertices*

$$H = - \sum_{\square} \sigma_i \sigma_j \sigma_k \sigma_l$$

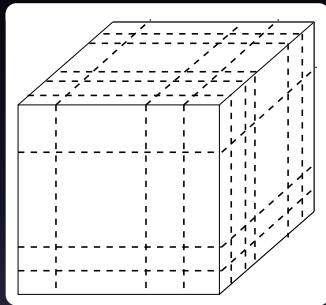
Not edges

$$H = - \sum_{\square} U_{ij} U_{jk} U_{kl} U_{li}$$

Plaquette Hamiltonian Groundstates: Single cube



Ground states, Low T: Lattice



Persists into low temperature phase (Wegner, Pietig)

Degeneracy 2^{3L}

Consequences: FSS

First order FSS: Boundary Conditions

Pirogov-Sinai Theory (Borgs/Kotecký)

$$Z(\beta) = \left[e^{-\beta L^d f_d} + q e^{-\beta L^d f_o} \right] [1 + \dots]$$

Fixed boundaries ($1/L$ leading term)

$$Z(\beta) = \left[e^{-\beta(L^d f_d + L^{d-1} \hat{f}_o)} + q e^{-\beta(L^d f_o + L^{d-1} \hat{f}_d)} \right]$$

First order FSS: Exponential Degeneracy

Exponential degeneracy

$$Z(\beta) = \left[e^{-\beta L^d f_d} + 2^{3L} e^{-\beta L^d f_o} \right]$$

$1/L^2$ in $d = 3$

$$Z(\beta) = \left[e^{-\beta L^d f_d} + e^{(3 \ln 2)L} e^{-\beta L^d f_o} \right]$$

FSS: Specific Heat

Probability of being in any of the states

$$p_o \propto e^{-\beta L^d \hat{f}_o} \text{ and } p_d \propto e^{-\beta L^d \hat{f}_d}$$

Time spent in the ordered states $\propto qp_o$

$$W_o/W_d \simeq q e^{-L^d \beta \hat{f}_o} / e^{-\beta L^d \hat{f}_d}$$

Expand around β^∞

$$0 = \ln q + L^d \Delta \hat{e} (\beta - \beta^\infty) + \dots$$

Solve for specific heat peak

$$\beta^{C_v^{\max}}(L) = \beta^\infty - \frac{\ln q}{L^d \Delta \hat{e}} + \dots$$

FSS: Specific Heat, degenerate case

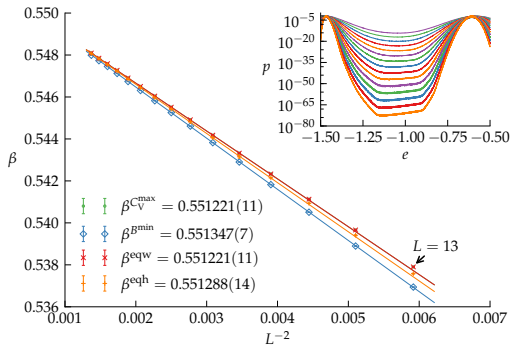
With $q \propto 2^{3L} = e^{(3 \ln 2)L}$

$$\beta^{C_V^{\max}}(L) = \beta^\infty - \frac{\ln q}{L^d \Delta \hat{e}} + \dots$$

become

$$\beta^{C_V^{\max}}(L) = \beta^\infty - \frac{3 \ln 2}{L^{d-1} \Delta \hat{e}} + \dots$$

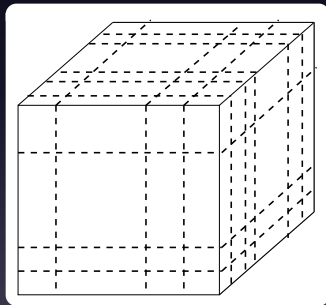
Plaquette Hamiltonian fits



Consequences: Order parameter

Order Parameter

Magnetization won't do
Not a gauge theory



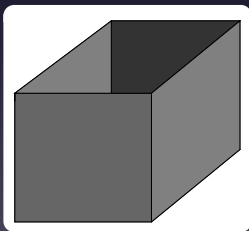
Anisotropic (“Fuki-Nuke”) Model

Consider anisotropic variant

Suzuki 1973, Jonsson/Savvidy 2000, Castelnovo *et.al.* 2010

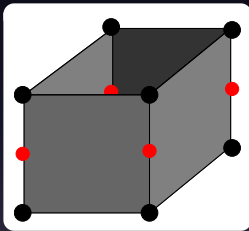
$$H_{\text{aniso}} = -J_x \sum_{\square yz} \sigma_i \sigma_j \sigma_k \sigma_l - J_y \sum_{\square xz} \sigma_i \sigma_j \sigma_k \sigma_l$$

Set $J_z = 0$, $J_x = J_y = 1$. No horizontal plaquettes



Anisotropic Model

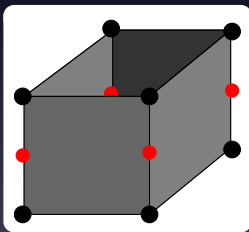
Define new spins τ from pairs of σ 's



$$\tau_{x,y,z} = \sigma_{x,y,z} \sigma_{x,y,z+1}$$

Anisotropic Model = Stack of $2D$ Ising

$$H_{\text{fuki-nuke}} = - \sum_{x=1}^L \sum_{y=1}^L \sum_{z=1}^L (\tau_{x,y,z} \tau_{x+1,y,z} + \tau_{x,y,z} \tau_{x,y+1,z}) ,$$



Anisotropic Model Order Parameter

Single layer

$$m_{2d,z} = \left\langle \frac{1}{L^2} \sum_{x=1}^L \sum_{y=1}^L \tau_{x,y,z} \right\rangle$$

In terms of original variables

$$m_{2d,z} = \left\langle \frac{1}{L^2} \sum_{x=1}^L \sum_{y=1}^L \sigma_{x,y,z} \sigma_{x,y,z+1} \right\rangle$$

Add 'em up

$$m_{\text{abs}} = \left\langle \frac{1}{L^3} \sum_{z=1}^L \left| \sum_{x=1}^L \sum_{y=1}^L \sigma_{x,y,z} \sigma_{x,y,z+1} \right| \right\rangle ,$$

Isotropic case

Hypothesis: Same order parameter works

Hashizume and Suzuki (2011)

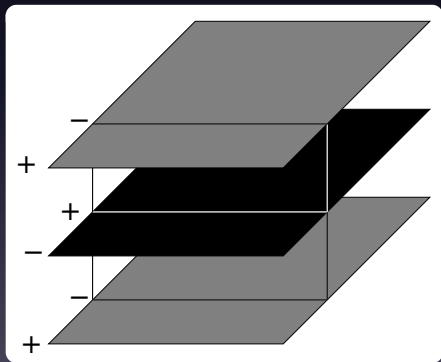
Take layers, add 'em up

$$m_{\text{abs}}^z = \left\langle \frac{1}{L^3} \sum_{z=1}^L \left| \sum_{x=1}^L \sum_{y=1}^L \sigma_{x,y,z} \sigma_{x,y,z+1} \right| \right\rangle$$

$$m_{\text{sq}}^z = \left\langle \frac{1}{L^5} \sum_{z=1}^L \left(\sum_{x=1}^L \sum_{y=1}^L \sigma_{x,y,z} \sigma_{x,y,z+1} \right)^2 \right\rangle$$

Isotropic case: Effect of flips

$$m_{\text{abs}}^z = \left\langle \frac{1}{L^3} \sum_{z=1}^L \left| \sum_{x=1}^L \sum_{y=1}^L \sigma_{x,y,z} \sigma_{x,y,z+1} \right| \right\rangle$$



Isotropic case: numerical investigation

Strong first order PT

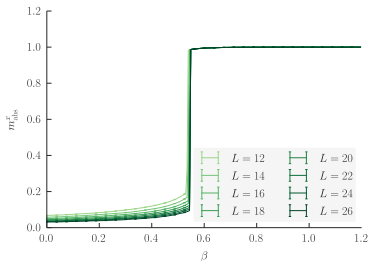
Multicanonical simulation

Correct order parameter: predicts PT point? isotropic?

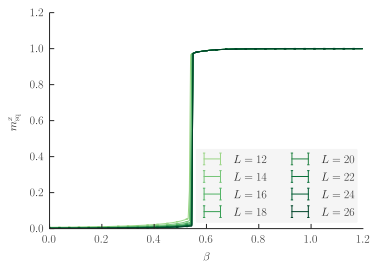
FSS properties?

Numerical results: order parameters

Order parameter m_{abs}^x



Order parameter m_{sq}^x



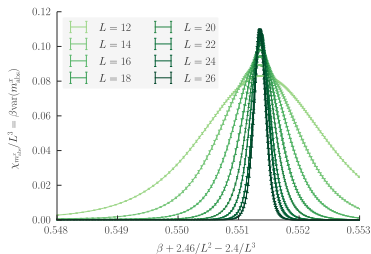
m^y and m^z identical to m^x - isotropy restored

Numerical results: scaled susceptibilities

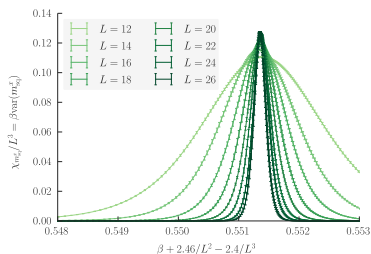
$$\chi(\beta) = \beta L^3 (\langle m^2 \rangle(\beta) - \langle m \rangle(\beta)^2)$$

$$\beta \chi_{m_{\text{abs}}}^x(L) = 0.551\,37(3) - 2.46(3)/L^2 + 2.4(3)/L^3$$

Susceptibility for m_{abs}^x



Susceptibility for m_{sq}^x



Conclusions

Exponential degeneracy gives $1/L^2$ corrections in 3D

Suzuki was right - isotropic model displays Fuki-Nuke order

OP Scaling agrees well with energetic observables

Perspectives

Cases where *ground state* is exponentially degenerate but low-T phase is not: AFM Ising on FCC

Scaling at first order Quantum PTs

References

G. K. Savvidy and F.J. Wegner, Nucl. Phys. B **413** (1994) 605.

Y. Hashizume and M. Suzuki, Int. J. Mod. Phys. B **25** (2011) 73; Int. J. Mod. Phys. B **25** (2011) 3529.

M. Mueller, W. Janke and D. A. Johnston, Phys. Rev. Lett. **112** (2014) 200601.

M. Mueller, D. A. Johnston and W. Janke, Nucl. Phys. B **888** (2014) 214.

M. Mueller, W. Janke and D. A. Johnston, Nucl. Phys. B **894** (2015) 1.