Surface defects in 4d supersymmetric field theories via integrable lattice model

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Introduction

In this talk, we consider half-BPS surface defects in 4d supersymmetric theories, supported on 2d surface in 4d spacetime.

Physically, these are useful to determine the phase of the theories, like a Wilson and 't Hooft loops, and can exist as a BPS object in N=1,2,3,4 theories in 4d.

Mathematical-physically, often these defects are related with some "operation" in integrable models/CFTs behind 4d supersymmetric theories.

N=2 class S theories

Let us focus on the theories in so-called class S, which are obtained by compactifying M5-branes (6d (2,0) theory) on a Riemann surface.

[Gaiotto, Gaiotto-Moore-Neitzke] [Pestun's and Teschner's talks]



The instanton (omega-deformed) partition function of this theory is related to the conformal block of 2d Virasoro or W algebra on this Riemann surface. [Alday-Gaiotto-Tachikawa]

Surface defects in class S

Half-BPS surface defects are described in 2d CFT:

[Alday-Gaiotto-Gukov-Tachikawa-Verlinde, Alday-Tachikawa....]

	0		2	3	4	5	6	7	8	9	
M5 (6d (2,0))	X	X	X	X			Х				×
M2 (codim 4)	Х	×						Х			
M5' (codim 2)	X	X			X		Х	X			X



Index of N=2 theories

We here consider **the supersymmetric index**, or the partition function on S¹xS³. This quantity is a bit simpler than the previous instanton partition function, but has a rich connection to the integrable model. [RomeIsberger, Kinney-Maldacena-Minwalla-Raju,...]



M5'-brane defect



???

Brane tiling model

To tackle this, we argue another class of N=1 theories, Brane tiling models which have some overlap with N=2 class S theories:

[Franco-Hanany-Vegh-Wecht-Kennaway]

This is realized on intersecting D5-NS5 system in Type IIB (or D3-branes on toric CY), and generally is 4d N=1 theory whose matter content is specified by a quiver diagram, like



Brane tiling/integrable lattice model

These N=1 models are known to be related to integrable lattice model on the same 2d surface on which the tiling is defined **[Spiridonov, Yamazaki]**

S³xS¹ partition function (or superconformal index) of Brane tiling model

Partition function of integrable lattice model

by [Bazhanov-Sergeev]



A model on a lattice



- to a white face, we associate continuous spin variables z_i
- a crossing of lines gives the interaction of spins

Integrable lattice model

[Bazhanov-Sergeev]

They show that the following star-triangle relation is solved



by assigning a particular Boltzmann weight to a interaction and selfinteraction, written by elliptic gamma function:

$$\Gamma(a; p, q) = \prod_{i,j>0} \frac{1 - a^{-1} p^{i+1} q^{j+1}}{1 - a p^i q^j}$$

This leads to integrability of the model.

Surface defects

Actually one can introduce **a different kind of rapidity line** in integrable model. From the intersection of this and the ordinary rapidity line is called L-operator (or Lax operator)

$$\check{L}_{ij} = i \cdots \downarrow j$$

<u>Proposal</u>: a class of surface defects is represented by the transfer matrix of L-operators in the corresponding integrable lattice model.

$$c$$
 z_1 z_2 \cdots z_n $z_$

Surface defects in class S

We consider the special case of the brane-tiling model where the supersymmetry in 4d is N=2, then introduce the surface defects. This is associated with L-operator:



Brane tiling and integrable model Surface defect as transfer matrix Surface defects in class S theories

Brane tiling and integrable model

Type IIB brane configuration

Let us consider the following brane configuration in Type IIB:

	0		2	3	4	5	6	7	8	9
N D5	×	X	X	Х	×		X			
NS5	X	Х	Х	Х	Х	Х				
NS5'	×	Х	Х	Х			X	X		

which preserves 4 supercharges (N=1 in 4d).

Let x^4 and x^6 coordinates parametrize T². Then, NS5-branes intersect with D5-branes **along curves in x^4-x^6 plane**.



Brane intersection

At the intersection of D5-brane and NS5-brane we can resolve as



We call the lines representing NS5-branes along half-line as rapidity lines.

In general, the diagram from rapidity lines on T² or other surface is called as **brane tiling**.

To a brane tiling without |q| > 1 region there is an associated an N=1 4d quiver gauge theory which is a worldvolume theory on D5-branes:

- to a face with q=0, we associate an SU(N) gauge group with fugacity z_i.
- a crossing of rapidity lines, each of which has continuous parameter, gives rise to a bi-fundamental chiral multiplet



Let us take a pair of rapidity lines. By using this doubled notation we can avoid the region with |q| > 1

$$(a, b) \longrightarrow = a \xrightarrow{a \longrightarrow b} b$$

A quiver gauge theory

We choose the (N,0) background, which leads to the following:



For example, this produces the following quiver gauge theory





R-matrix and Yang-Baxter eq

We associate an R-matrix to the crossing of lines:

Index of
$$\left(\begin{array}{c} (a_i, b_j) & & \\ \hline & & \\ (a_i, b_i) & & \\ \hline & & \\ (a_j, b_j) \end{array} \right) = R_{ij}^{\diamond}(a_i, b_i; a_j, b_j)_{z_i, z_j}^{w_i, w_j}$$

This defines a map $R_{ij}^{\diamond}: V_i^{\diamond} \otimes V_j^{\diamond} \to V_j^{\diamond} \otimes V_i^{\diamond}$, where V_i^{\diamond} is the space of symmetric meromorphic functions of z_i or w_i .

The R-matrix satisfies the Yang-Baxter equation

 $R_{12}^{\diamond}(p_1, p_2)R_{13}^{\diamond}(p_1, p_3)R_{23}^{\diamond}(p_2, p_3) = R_{23}^{\diamond}(p_2, p_3)R_{13}^{\diamond}(p_1, p_3)R_{12}^{\diamond}(p_1, p_2)$



$$p_i = (a_i, b_i)$$

This is indeed satisfied because of the brane construction after uplift to M-theory.

Surface defect

as transfer matrix

Half-BPS surface defect

We consider **D3-brane insertion** in the previous Type IIB brane configuration which preserves 2 supercharges:

	0		2	3	4	5	6	7	8	9
N D5	×	×	×	×	×		×			
NS5	×	×	×	×	×	×				
NS5'	×	X	X	X			×	×		
D3	×	Х			Х			X		

Depending on the NS5-brane worldvolume direction, we have, for example:



These can be regarded as **symmetric** and **antisymmetric** representation of SU(N).

"L-operator"

We denote the projection of D3-branes in x^4-x^6 surface by the dashed arrow:

C

Similar to the line from NS5, the dashed line associates the fugacity c, and the representation space W_R of rep. R.

By using this and zigzag path, we now have the following three types of intersection:

$$\check{R}_{ij} = i \xrightarrow{j}, \quad \check{L}_{ij} = i \xrightarrow{j}, \quad \check{R}_{ij} = i \xrightarrow{j}_{C_1}, \quad \check{R}_{ij} = i \xrightarrow{j}_{C_2}.$$

 $L(c, (a_j, b_j)): W_i \otimes V_j \to V_j \otimes W_i$

"Yang-Baxter equations"

Accordingly, three types of Yang-Baxter like equations can be considered:

One dashed line:



Two dashed lines:



Three dashed lines:



Again these are indeed satisfied because of the brane construction after uplift to M-theory.

L-operator (N=2 case)

Let us focus on N=2 (SU(2)) case and the representation of surface defect is the fundamental one. It is known that there is an operator in integrable model, so-called L-operator L: $W_i \otimes V_j \rightarrow V_j \otimes W_i$

This can be represented by $2x^2$ matrix whose entries are operators acting on function in V_i.

Yang-Baxter equations:



Proposal

[Yagi-KM]

Sklyanin found that the following operator solves the RLL relation:

$$L = \sum_{a=0}^{3} w_a \ \sigma_a \otimes \mathbf{S}_a$$

where S_a satisfies the so-called Sklyanin algebra and is represented by a difference operator.

We propose that the transfer matrix from these L-operators is identified with the index of the surface defect

$$\operatorname{Tr}(L(c,(a_n,b_n))\dots L(c,(a_1,b_1))) = \sum_{s_1=\{\pm 1\}} \dots \sum_{s_n=\{\pm 1\}} \prod_{i=1}^n l(z_i^{s_i}, z_{i+1}^{s_{i+1}}; \frac{b_i}{c}, \frac{a_{i+1}}{c}) \prod_{j=1}^n \Delta_j^{s_j/2}.$$

where

$$(\Delta_i^{\pm 1/2} f)(z_i) = f(q^{\pm 1/2} z_i) \qquad l(z_i, z_{i+1}; \frac{b_i}{c}, \frac{a_{i+1}}{c}) = \frac{1}{\theta(z_i^2)} \theta(\sqrt{pq} \frac{c^2}{b_i a_{i+1}} \frac{z_{i+1}}{z_i}) \theta(\sqrt{pq} \frac{b_i}{a_{i+1}} z_i z_{i+1}).$$

Generalizations

- We can consider the surface defects supported on S¹xS¹. The S¹ can be one of two circles in S³ (|ξ₁|² + |ξ₂|² = 1). Having two choices corresponds to modular double of Sklyanin algebra. (p↔q)
- General N and general representations have been studied in [Yagi, '17]

Surface defect in class S theory

Class S theories

Now we come back to class S theory. Actually some theories in class S are a special case of brane tiling models.

Let us consider the following tiling:

identify upper nodes with lower nodes





M5 on cylinder

The previous brane tiling is produced by the brane system like:

	0		2	3	4	5	6	7	8	9
N D5	×	X	X	×	X		×			
NS5	X	Х	Х	Х	Х	Х				



	0		2	3	4	5	6	7	8	9
N D4	×	×	×	×			×			
NS5	X	Х	X	Х	Х	Х				





A surface defect

We add to this a D3 brane to describe a surface defect.

Depending on the directions, we have two types of defects:

		0		2	3	4	5	6	7	8	9
Ι.	N D5	×	Х	Х	Х	Х		×			
	NS5	×	X	×	X	X	×				
	D 3	Х	Х			Х			Х		



	0		2	3	4	5	6	7	8	9
N D5	×	×	×	×	×		\times			
NS5	×	×	×	×	×	×				
D3'	X	×					×	×		

2.



The first case:

From the proposal from integrable lattice model, the inclusion of the surface defect corresponds just to the trace of single L-operator:



The difference operator acts on the flavor or gauge fugacity z

	0		2	3	4	5	6	7	8	9
N D 4	×	×	×	×			×			
NS5	×	×	×	×	X	×				
D2	×	×						×		

N M5-branes on



[Gaiotto-Rastelli-Razamat]

The second case:

This corresponds to the trace of multiple L-operators:



	0		2	3	4	5	6	7	8	9
N D4	×	Х	Х	X	Х		X			
NS5	×	\times	\times	\times	×	×				
D 4	×	×			×		×	×		

N M5-branes + M5 on



Conclusion

We propose that the index of the half-BPS surface defect from D3branes in brane tiling model can be identified with the transfer matrix of L-operators in the corresponding integrable model.

This gives the unified picture of the indices of the M2- and M5branes defects in class S theories.

Outlook

M2-M5' defect from Baxter's R matrix
another T-dual gives line defects in class S
3d reduction