Multiscale modelling, analysis and simulations of plant biomechanics and cellular signalling processes

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Plant cell walls: microstructure, mechanics & chemistry

Microscopic structure of plant cell walls

- cellulose microfibrils
- cell wall matrix of pectin, hemicellulose, water, enzymes
- allows for anisotropic cell expansion

Interactions between mechanics and chemistry

- mechanical forces can break load-bearing cross-links
- dynamics of cross-links influences mechanical properties of plant cell wall matrix, of plant cell wall matrix, by cell wall pectins



B pH+

activity of cell wall remodeling enzymes



Mechanics (hyperelastic material)

div
$$\mathbf{T} = 0$$
, $\mathbf{T} = J_e^{-1} \mathbf{F}_e \frac{\partial W(\mathbf{F}_e)}{\partial \mathbf{F}_e}$

+ boundary conditions

 $\mathbf{F} = \mathbf{I} + \nabla \mathbf{u}$ deformation gradient $\mathbf{F} = \mathbf{F}_e \mathbf{F}_g$ decomposition in elastic & growth $J_e = \det(\mathbf{F}_e), J_g = \det(\mathbf{F}_g)$

Ogden, *Nonlinear elastic deform*. 1984 Rodriguez, Hoger, McCulloch, *J Biomech*. 1994 Goriely, Moulton, Vandiver, *EPL* 2010 Goriely, Ben Amar, *J Mech.Phys.Solids* 2005 Goriely, Ben Amar, *Biomech.Model.Mechan*. 2007 Huang, Becker, Jones, *J Mech.Phys.Solids* 2012

or $\mathbf{T} = \mathbf{F}_e \frac{\partial W(\mathbf{F}_e)}{\partial \mathbf{F}} - p\mathbf{I}$

Calcium-pectin chemistry

- methylestrified pectin: b₁
- ► demethylestrified pectin: *b*₂
- pectin-calcium cross links: b₃
- enzyme PME: p
- calcium ions: c

$$\partial_t n - \operatorname{div}(D_n \nabla n) = g(n, \nabla \mathbf{u})$$

$$n \in \{b_1, b_2, b_3, p, c\}$$

$$\bigwedge^{\mathsf{PME}} \longrightarrow^{\mathsf{PME}} + \bullet^{\mathsf{+}} \mathsf{H}^{\mathsf{+}}$$

Plant cell walls: mechanics & chemistry

Linear elasticity or viscoelasticity

div
$$\mathbf{T} = 0$$
 in G , $\mathbf{T} \cdot \mathbf{n} = -P\mathbf{n}$ on ∂G
 $\mathbf{T} = \left(\mathbb{E}_{M}(b_{3})\chi_{G_{M}} + \mathbb{E}_{F}\chi_{G_{F}}\right)\mathbf{e}(\mathbf{u}_{e})$
 $\mathbf{T} = \left(\mathbb{E}_{M}(b_{3})\mathbf{e}(\mathbf{u}_{e}) + \mathbb{V}_{M}(b_{3})\mathbf{e}(\partial_{t}\mathbf{u}_{e})\right)\chi_{G_{M}} + \mathbb{E}_{F}\chi_{G_{F}}\mathbf{e}(\mathbf{u}_{e})$
 $\mathbf{e}(\mathbf{u}_{e})_{ij} = \frac{1}{2}(\partial_{x_{i}}\mathbf{u}_{e,j} + \partial_{x_{j}}\mathbf{u}_{e,i})$

Reaction-diffusion equations for chemical reactions

$$\partial_t n - \operatorname{div}(D_n \nabla n) = g_n(n, \mathcal{R}(\mathbf{e}(\mathbf{u}_e))),$$

- demethyl-esterified pectin can decay
- formation and destruction of calcium-pectin cross-links

$$\mathcal{R}(\mathbf{e}(\mathbf{u}_e)) = \left(\operatorname{tr} \left(\mathbb{E}_M(b_3) \chi_{G_M} + \mathbb{E}_F \chi_{G_F} \right) \mathbf{e}(\mathbf{u}_e) \right)^+$$

$$rac{\kappa}{1+eta b_2} b_1 \, p$$

 $n \in \{b_1, b_2, b_3, p, c\}$

$$-2g(c)b_2+2\kappa b_3 \mathcal{R}(\mathbf{e}(\mathbf{u}_e))$$

or
$$\mathcal{R}(\mathbf{e}(\mathbf{u}_e)) = (\operatorname{tr} \mathbf{e}(\mathbf{u}_e))^+$$

MP, B. Seguin, ESAIM M2AN, 2016



Microscopic Model In $(0, T) \times G$

 $\operatorname{div}(\mathbb{E}^{\varepsilon}(\boldsymbol{b}_{3}^{\varepsilon}, x) \mathbf{e}(\mathbf{u}_{e}^{\varepsilon})) = \mathbf{0}$

or

 $\operatorname{div}(\mathbb{E}^{\varepsilon}(\boldsymbol{b}_{3}^{\varepsilon}, \boldsymbol{x})\mathbf{e}(\mathbf{u}_{e}^{\varepsilon}) + \mathbb{V}^{\varepsilon}(\boldsymbol{b}_{3}^{\varepsilon}, \boldsymbol{x})\mathbf{e}(\partial_{t}\mathbf{u}_{e}^{\varepsilon})) = \mathbf{0}$

 $\mathbb{E}^{\varepsilon}(\xi, x) = \mathbb{E}(\xi, \hat{x}/\varepsilon), \quad \mathbb{V}^{\varepsilon}(\xi, x) = \mathbb{V}(\xi, \hat{x}/\varepsilon), \text{ where }$

 $\mathbb{E}(\xi, \hat{y}) = \mathbb{E}_{M}(\xi) \chi_{\hat{Y}_{M}}(\hat{y}) + \mathbb{E}_{F} \chi_{\hat{Y}_{F}}(\hat{y}), \quad \mathbb{V}(\xi, \hat{y}) = \mathbb{V}_{M}(\xi) \chi_{\hat{Y}_{M}}(\hat{y}) \quad \text{are } \hat{Y} - \text{periodic}, \\ \hat{Y} = Y \cap \{x_{3} = \text{const}\}$

In $(0, T) \times G_M^{\varepsilon}$

 $\begin{aligned} \partial_{t} p^{\varepsilon} &= \operatorname{div}(D_{p} \nabla p^{\varepsilon}) \\ \partial_{t} b_{1}^{\varepsilon} &= \operatorname{div}(D_{b_{1}} \nabla b_{1}^{\varepsilon}) - f(b_{1}^{\varepsilon}, b_{2}^{\varepsilon}, p^{\varepsilon}) \\ \partial_{t} b_{2}^{\varepsilon} &= \operatorname{div}(D_{b_{2}} \nabla b_{2}^{\varepsilon}) + f(b_{1}^{\varepsilon}, b_{2}^{\varepsilon}, p^{\varepsilon}) - 2g(c^{\varepsilon})b_{2}^{\varepsilon} + 2\kappa b_{3}^{\varepsilon} \mathcal{R}(\mathbf{e}(\mathbf{u}_{e}^{\varepsilon})) \\ \partial_{t} c^{\varepsilon} &= \operatorname{div}(D_{c} \nabla c^{\varepsilon}) - g(c^{\varepsilon})b_{2}^{\varepsilon} + \kappa b_{3}^{\varepsilon} \mathcal{R}(\mathbf{e}(\mathbf{u}_{e}^{\varepsilon})) \\ \partial_{t} b_{3}^{\varepsilon} &= \operatorname{div}(D_{b_{3}} \nabla b_{3}^{\varepsilon}) + g(c^{\varepsilon})b_{2}^{\varepsilon} - \kappa b_{3}^{\varepsilon} \mathcal{R}(\mathbf{e}(\mathbf{u}_{e}^{\varepsilon})) \end{aligned}$



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Existence of solutions of the model for the cell wall

- For $b^{\varepsilon} \in L^{\infty}(0, T; L^{\infty}(G_{M}^{\varepsilon}))$ with $b^{\varepsilon} \geq 0$
- $\exists \mathbf{u}_{e}^{\varepsilon} \in L^{\infty}(0, T; \mathcal{W}(G)) \text{ satisfying elasticity or viscoelasticity problem} \\ \|\mathbf{e}(\mathbf{u}_{e}^{\varepsilon,1} \mathbf{u}_{e}^{\varepsilon,2})\|_{L^{\infty}(0,T;L^{2}(G))}^{2} \leq C \|b^{\varepsilon,1} b^{\varepsilon,2}\|_{L^{\infty}(0,T;L^{\infty}(G_{M}^{\varepsilon}))}^{2}$

• For
$$\mathbf{u}_e^{\varepsilon} \in L^{\infty}(0, T; \mathcal{W}(G))$$
 such that

 $\|\mathbf{u}_{e}^{\varepsilon}\|_{L^{\infty}(0,T;\mathcal{W}(G))}\leq C,$

 \exists non-negative unique weak solution $(b^{\varepsilon}, c^{\varepsilon})$ such that

$$\|\boldsymbol{b}^{\varepsilon,1}-\boldsymbol{b}^{\varepsilon,2}\|_{L^{\infty}(0,\tilde{T};L^{\infty}(G_{M}^{\varepsilon}))}^{2} \leq C\tilde{T} \|\boldsymbol{e}(\boldsymbol{u}_{e}^{\varepsilon,1}-\boldsymbol{u}_{e}^{\varepsilon,2})\|_{L^{\infty}(0,\tilde{T};L^{2}(G))}^{2}, \quad \tilde{T} \in (0,T]$$

• $\mathcal{K}: L^{\infty}(G_{M,\tilde{T}}^{\varepsilon}) \to L^{\infty}(G_{M,\tilde{T}}^{\varepsilon})$ by $\mathcal{K}(\tilde{b}^{\varepsilon}) = b^{\varepsilon}$

$$\mathcal{W}(G) = \left\{ \mathbf{u} \in H^1(G; \mathbb{IR}^3) \mid \int_G \mathbf{u} \, dx = \mathbf{0}, \int_G [(\nabla \mathbf{u})_{12} - (\nabla \mathbf{u})_{21}] dx = \mathbf{0}, \ \mathbf{u} \text{ periodic in } x_3 \right\}$$

Existence of solutions of the model for the cell wall

• For
$$b^{\varepsilon} \in L^{\infty}(0, T; L^{\infty}(G_{M}^{\varepsilon}))$$
 with $b^{\varepsilon} \geq 0$

 $\exists \mathbf{u}_e^{\varepsilon} \in L^{\infty}(0, T; \mathcal{W}(G))$ satisfying elasticity or viscoelasticity problem

$$\|\mathbf{e}(\mathbf{u}_{e}^{\varepsilon,1}-\mathbf{u}_{e}^{\varepsilon,2})\|_{L^{\infty}(0,T;L^{2}(G))}^{2} \leq C\|b^{\varepsilon,1}-b^{\varepsilon,2}\|_{L^{\infty}(0,T;L^{\infty}(G_{M}^{\varepsilon}))}^{2}$$

• For
$$\mathbf{u}_e^{\varepsilon} \in L^{\infty}(0, T; \mathcal{W}(G))$$
 such that

$$\|\mathbf{u}_{e}^{\varepsilon}\|_{L^{\infty}(0,T;\mathcal{W}(G))}\leq C,$$

 \exists non-negative unique weak solution $(b^{\varepsilon}, c^{\varepsilon})$ such that

$$\|\boldsymbol{b}^{\varepsilon,1}-\boldsymbol{b}^{\varepsilon,2}\|_{L^{\infty}(0,\tilde{T};L^{\infty}(G_{M}^{\varepsilon}))}^{2} \leq C\tilde{T} \|\mathbf{e}(\mathbf{u}_{e}^{\varepsilon,1}-\mathbf{u}_{e}^{\varepsilon,2})\|_{L^{\infty}(0,\tilde{T};L^{2}(G))}^{2}, \quad \tilde{T} \in (0,T]$$

•
$$\mathcal{K}: L^{\infty}(G^{\varepsilon}_{\mathcal{M},\tilde{T}}) \to L^{\infty}(G^{\varepsilon}_{\mathcal{M},\tilde{T}})$$
 by $\mathcal{K}(\tilde{b}^{\varepsilon}) = b^{\varepsilon}$

Main tools: Gagliardo-Nirenberg inequality & Moser-Alikakos iteration
 technique

Multiscale analysis

Aim of multiscale analysis:

to defined macroscopic behaviour of a biological or physical system by taking microscopic processes and microstructure into account



- H-, I-, and G- convergences
- periodic and locally-periodic unfolding operators

Two-scale convergence

• A special type of convergence in L^p , 1 and <math>1/p + 1/q = 1

Definition. $\{u^{\varepsilon}\} \subset L^{p}(\Omega)$ two-scale converge to $u, u \in L^{p}(\Omega \times Y)$ iff for any $\phi \in L^{q}(\Omega, C_{per}(Y))$

$$\lim_{\varepsilon \to 0} \int_{\Omega} u^{\varepsilon}(x) \phi\left(x, \frac{x}{\varepsilon}\right) dx = \int_{\Omega} f_{Y} u(x, y) \phi(x, y) dx dy.$$

Notice:

$$u^{\varepsilon} \rightharpoonup \int_{Y} u(\cdot, y) dy$$
 weakly in $L^{p}(\Omega)$



Definition. $\{u^{\varepsilon}\} \subset L^2(\Gamma^{\epsilon})$ two-scale converge to $u, u \in L^2(\Omega \times \Gamma)$ iff for $\psi \in C(\overline{\Omega}, C_{per}(Y))$:

$$\lim_{\varepsilon \to 0} \varepsilon \int_{\Gamma^{\varepsilon}} u^{\epsilon}(x) \psi(x, x/\varepsilon) d\gamma_{x} = \frac{1}{|Y|} \int_{\Omega} \int_{\Gamma} u(x, y) \psi(x, y) dx d\gamma_{y}$$

Y

Convergence results

 $b, c \in L^2(0, T; H^1(G)), \quad c^1, b^1 \in L^2(G_T; H^1_{per}(\hat{Y}))$ $\mathbf{u}_e \in L^\infty(0, T; \mathcal{W}(G)), \quad \mathbf{u}_e^1 \in L^2(G_T; H^1_{per}(\hat{Y}))$

$$b^{\varepsilon}
ightarrow b, \quad c^{\varepsilon}
ightarrow c$$
 weakly in $L^{2}(0, T; H^{1}(G))$
 $\nabla b^{\varepsilon}
ightarrow \nabla b + \hat{\nabla}_{y} b^{1}, \quad \nabla c^{\varepsilon}
ightarrow \nabla c + \hat{\nabla}_{y} c^{1}$ weakly two-scale
 $b^{\varepsilon}
ightarrow b, \quad c^{\varepsilon}
ightarrow c$ strongly in $L^{2}(G_{T})$

$$\begin{aligned} \mathbf{u}_{e}^{\varepsilon} \rightharpoonup \mathbf{u}_{e} & \text{weakly}^{*} \text{ in } L^{\infty}(0, T; \mathcal{W}(G)) \\ \nabla \mathbf{u}_{e}^{\varepsilon} \rightharpoonup \nabla \mathbf{u}_{e} + \hat{\nabla}_{y} \mathbf{u}_{e}^{1} & \text{weakly two-scale} \end{aligned}$$

 $\mathbf{e}(\mathbf{u}_e^{\varepsilon}) \rightarrow \mathbf{e}(\mathbf{u}_e) + \mathbf{e}_y(\mathbf{u}_e^1)$ strongly two-scale

 $\hat{Y} = Y \cap \{x_3 = \text{const}\}, \qquad G_T = (0, T) \times G$

Microscopic Model

In $(0, T) \times G$

 $\operatorname{div}(\mathbb{E}^{\varepsilon}(\mathbf{b}^{\varepsilon}, x)\mathbf{e}(\mathbf{u}_{e}^{\varepsilon})) = \mathbf{0}$

or

 $\operatorname{div}(\mathbb{E}^{\varepsilon}(\boldsymbol{b}^{\varepsilon}, \boldsymbol{x}) \mathbf{e}(\mathbf{u}_{e}^{\varepsilon}) + \mathbb{V}^{\varepsilon}(\boldsymbol{b}^{\varepsilon}, \boldsymbol{x}) \mathbf{e}(\partial_{t}\mathbf{u}_{e}^{\varepsilon})) = \mathbf{0}$



In $(0, T) \times G_M^{\varepsilon}$

 $\partial_t b^{\varepsilon} = \operatorname{div}(D_b \nabla b^{\varepsilon}) + g_b(b^{\varepsilon}, c^{\varepsilon}, \mathcal{R}(\mathbf{e}(\mathbf{u}_e^{\varepsilon})))$ $\partial_t c^{\varepsilon} = \operatorname{div}(D_c \nabla c^{\varepsilon}) + g_c(b^{\varepsilon}, c^{\varepsilon}, \mathcal{R}(\mathbf{e}(\mathbf{u}_e^{\varepsilon})))$



or $\mathcal{R}(\mathbf{e}(\mathbf{u}_e^{\varepsilon})) = (\operatorname{tr} \mathbf{e}(\mathbf{u}_e^{\varepsilon}))^+$

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$$\mathcal{R}(\mathbf{e}(\mathbf{u}_{e}^{\varepsilon})) = \left(\operatorname{tr}\left(\mathbb{E}_{M}(\mathbf{b}^{\varepsilon})\chi_{G_{M}} + \mathbb{E}_{F}\chi_{G_{F}}\right)\mathbf{e}(\mathbf{u}_{e}^{\varepsilon}) \right)^{+}$$

Numerical simulations for plant cell wall model

Macroscopic model for plant cell wall biomechanics

$$\operatorname{div}(\mathbb{E}_{\operatorname{hom}}(\boldsymbol{b}_3)\,\mathbf{e}(\mathbf{u}_e))=\mathbf{0}\qquad \text{ in } G_T$$

$$\partial_t b = \operatorname{div}(\mathcal{D}_b \nabla b) + g_b(b, c, R(\mathbf{e}(\mathbf{u}_e))) \quad \text{in } G_T$$
$$\partial_t c = \operatorname{div}(\mathcal{D}_c \nabla c) + g_c(b, c, R(\mathbf{e}(\mathbf{u}_e))) \quad \text{in } G_T$$

$$R(\mathbf{e}(\mathbf{u}_{e})) = \left(\operatorname{tr}\left(\mathbb{E}_{\operatorname{hom}}(b_{3})\,\mathbf{e}(\mathbf{u}_{e})\right)\right)^{+} \text{ or } \left(\operatorname{tr}\mathbf{e}(\mathbf{u}_{e})\right)^{+}$$
$$\mathcal{D}_{\alpha,j3} = \mathcal{D}_{\alpha,3j} = \mathcal{D}_{\alpha}\delta_{3j}, \quad \mathcal{D}_{\alpha,ij} = \mathcal{D}_{\alpha} \oint_{\hat{Y}_{M}} \left[\delta_{ij} + \partial_{y_{j}}v_{\alpha}^{i}(y)\right] dy,$$
$$\alpha = b_{1}, b_{2}, b_{3}, c$$

MP, B. Seguin, *ESAIM M2AN*, 2016

$$\mathbb{E}_{\mathrm{hom},ijkl}(b_3) = \oint_{Y} \left[\mathbb{E}_{ijkl}(b_3, y) + \left(\mathbb{E}(b_3, y) \mathbf{e}_{y}(\mathbf{w}^{ij}) \right)_{kl} \right] dy$$

Macroscopic elasticity tensor

$$\mathbb{E}_{\text{hom},ijkl}(x,b_{3}) = \int_{Y} \left[\mathbb{E}_{Y,ijkl}(b_{3},y) + \mathbb{E}_{Y,ijpq}(b_{3},y)\mathbf{e}_{y}(\mathbf{w}^{kl})_{pq}(y) \right] dy$$

$$\operatorname{div}_{y} \left(\mathbb{E}_{Y}(b_{3},y)(\mathbf{e}_{y}(\mathbf{w}^{kl}) + \mathbf{b}^{kl}) \right) = \mathbf{0} \quad \text{in } Y$$

$$\int_{Y} \mathbf{w}^{kl} dy = \mathbf{0}, \qquad \mathbf{w}^{kl} \text{ is } Y \text{-periodic}$$

$$\mathbb{E}_{x}(b_{3},y) = \begin{cases} \mathbb{E}_{M}(b_{3}) & \text{if } y \in Y_{M} \\ \mathbb{E}_{F}(b_{3},y) = \begin{cases} \mathbb{E}_{M}(b_{3}) & \text{if } y \in Y_{M} \\ \mathbb{E}_{F}(b_{3},y) = \begin{cases} \mathbb{E}_{M}(b_{3}) & \text{if } y \in Y_{M} \\ \mathbb{E}_{F}(b_{3},y) = \begin{cases} \mathbb{E}_{M}(b_{3}) & \text{if } y \in Y_{M} \\ \mathbb{E}_{F}(b_{3},y) = \begin{cases} \mathbb{E}_{M}(b_{3}) & \text{if } y \in Y_{M} \\ \mathbb{E}_{F}(b_{3},y) = \begin{cases} \mathbb{E}_{M}(b_{3}) & \text{if } y \in Y_{M} \\ \mathbb{E}_{F}(b_{3},y) = \begin{cases} \mathbb{E}_{M}(b_{3}) & \text{if } y \in Y_{M} \\ \mathbb{E}_{F}(b_{3},y) = \begin{cases} \mathbb{E}_{M}(b_{3}) & \text{if } y \in Y_{M} \\ \mathbb{E}_{F}(b_{3},y) = \end{cases} \\ \mathbb{E}_{M}(b_{3}) & \text{if } y \in Y_{F}(b_{3},y) = \end{cases} \\ \mathbb{E}_{M}(b_{3}) & \text{if } y \in Y_{F}(b_{3},y) = \end{cases} \\ \mathbb{E}_{M}(b_{3}) & \text{if } y \in Y_{F}(b_{3},y) = \end{cases} \\ \mathbb{E}_{M}(b_{3}) & \text{if } y \in Y_{F}(b_{3},y) = \end{cases} \\ \mathbb{E}_{M}(b_{3}) & \text{if } y \in Y_{F}(b_{3},y) = \end{cases} \\ \mathbb{E}_{M}(b_{3}) & \text{if } y \in Y_{F}(b_{3},y) = \end{cases} \\ \mathbb{E}_{M}(b_{3}) & \text{if } y \in Y_{F}(b_{3},y) = \end{cases} \\ \mathbb{E}_{M}(b_{3}) & \text{if } y \in Y_{F}(b_{3},y) = \end{cases} \\ \mathbb{E}_{M}(b_{3}) & \text{if } y \in Y_{F}(b_{3},y) = \end{cases} \\ \mathbb{E}_{M}(b_{3}) & \text{if } y \in Y_{F}(b_{3},y) = \end{cases} \\ \mathbb{E}_{M}(b_{3}) & \text{if } y \in Y_{F}(b_{3},y) = \\ \mathbb{E}_{M}(b_{3}) & \text{if } y \in Y_{F}(b_{3},y) = \\ \mathbb{E}_{M}(b_{3}) & \text{if } y \in Y_{F}(b_{3},y) = \\ \mathbb{E}_{M}(b_{3}) & \text{if } y \in Y_{F}(b_{3},y) = \\ \mathbb{E}_{M}(b_{3}) & \text{if } y \in Y_{F}(b_{3},y) = \\ \mathbb{E}_{M}(b_{3}) & \text{if } y \in Y_{F}(b_{3},y) = \\ \mathbb{E}_{M}(b_{3}) & \text{if } y \in Y_{F}(b_{3},y) = \\ \mathbb{E}_{M}(b_{3}) & \text{if } y \in Y_{F}(b_{3},y) = \\ \mathbb{E}_{M}(b_{3}) & \text{if } y \in Y_{F}(b_{3},y) = \\ \mathbb{E}_{M}(b_{3}) & \text{if } y \in Y_{F}(b_{3},y) = \\ \mathbb{E}_{M}(b_{3}) & \text{if } y \in Y_{F}(b_{3},y) = \\ \mathbb{E}_{M}(b_{3}) & \text{if } y \in Y_{F}(b_{3},y) = \\ \mathbb{E}_{M}(b_{3}) & \text{if } y \in Y_{F}(b_{3},y) = \\ \mathbb{E}_{M}(b_{3}) & \text{if } y \in Y_{F}(b_{3},y) = \\ \mathbb{E}_{M}(b_{3}) & \text{if } y \in Y_{F}(b_{3},y) = \\ \mathbb{E}_{M}(b_{3}) & \text{if } y \in Y_{F}(b_{3},y) = \\ \mathbb{E}_{$$

Macroscopic elasticity tensor for cell wall

 $\mathbb{E}(y,b_3) = \mathbb{E}_M(b_3) \ \chi_{Y_M}(y) + \mathbb{E}_F \chi_{Y_F}(y)$

Cell wall matrix is assumed to be isotropic

 $\mathbb{E}_M(b_3) = E_M(b_3) \mathbb{E}_1 + \mathbb{E}_0 \longrightarrow \mathbb{E}_{\mathrm{hom}}(b_3) = E_M(b_3) \mathbb{E}_{\mathrm{hom},1} + \mathbb{E}_{\mathrm{hom},0}$

 $\mathbb{E}_M \mathbf{A} = 2\mu_M \mathbf{A} + \lambda_M (\operatorname{tr} \mathbf{A}) \mathbf{1}$ Lame moduli $\mu_M \quad \lambda_M$

$$E_M = \frac{\mu_M (2\mu_M + 3\lambda_M)}{\mu_M + \lambda_M}$$
 and $\nu_M = \frac{\lambda_M}{2(\mu_M + \lambda_M)}$

 $E_M(b_3) = 0.775 \ b_3 + 8.08 \ \text{MPa}$

(Zsivanovits, MacDougall, Smith, Ring Carbohydrate Research, 2004)

Microfibrils are transversally isotropic

| / | $\alpha_2 + \alpha_5$ | $\alpha_2 - \alpha_5$ | $lpha_3$ | 0 | 0 | $0 \rangle$ | (18393) | B 685 | 5 22277 | 0 | 0 | 0 \ |
|---|-----------------------|-----------------------|------------|------------|------------|-------------|---------------|--------|----------|-------|-------|-------|
| | $\alpha_2 - \alpha_5$ | $\alpha_2 + \alpha_5$ | $lpha_3$ | 0 | 0 | 0 | 6855 | 1839 | 3 22277 | 0 | 0 | 0 |
| | $lpha_3$ | $lpha_3$ | α_1 | 0 | 0 | 0 | 2227 | 7 2227 | 7 259901 | 0 | 0 | 0 |
| | 0 | 0 | 0 | α_4 | 0 | 0 | 0 | 0 | 0 | 84842 | 0 | 0 |
| | 0 | 0 | 0 | 0 | α_4 | 0 | 0 | 0 | 0 | 0 | 84842 | 0 |
| | 0 | 0 | 0 | 0 | 0 | α_5 | \setminus 0 | 0 | 0 | 0 | 0 | 5769/ |

 $E_F = 15,000 \text{ MPa}, \ \nu_{F1} = 0.3, \ n_F = 0.068, \ \nu_{F2} = 0.06, \ Z_F = 84,842 \text{ MPa}$

(Robert Moon, Review Wisconsin 2013-2014, Diddens et al., 2008)



nicrofibrils Ω_M^{ε} labeled. The surhidden) surface Γ_{ε} is facing the n the top and bottom of Ω_{ε} (b) Δ

Impact of the orientation of microfibrils



(BC1) Base case: $p = p_{\circ,1} = 0.209$ MPa and $f = 2.938p_{\circ,1}$ MPa

MP, B. Seguin, Bull Math Biology, 2016

Comparison with experimental results on tissue extension and compression

| RD | (BC1) parallel MF | (BC1) no shift | (BC1) 4 shifts | (BC1) rotated MF | (BC2) | (BC3) |
|------|-------------------|----------------|----------------|------------------|---------|---------|
| (C1) | 1.06497 | 1.06396 | 1.06984 | 1.00456 | 1.00683 | 1.00816 |
| (C2) | 1.06452 | 1.06391 | | 1.00462 | 1.00671 | 1.00823 |
| (C3) | 1.06234 | 1.06294 | 1.06680 | 1.00379 | 1.00672 | 1.00831 |
| (C4) | 1.06411 | 1.06320 | | 1.00497 | 1.00696 | 1.00811 |

Good agreement with experimental data on changes in inner or outer tissue length due to tissue tension elimination, Hejnowicz, Sievers, *J Exp.Botany* 1995:

relative displacement (RD) ranging between 0.38% and 6.98% versus experimental data: 0.3% - 4.99%

- (BC1) Base case: $p = p_{o,1} = 0.209$ MPa and $f = 2.938 p_{o,1}$ MPa
- (BC2) No tensile tractions: $p = p_{\circ,1}$ and f = 0
- (BC3) Different turgor pressures in neighbouring cells and no tensile tractions: $p_1 = p_4 = p_5 = p_8 = p_\circ$ and $p_2 = p_3 = p_6 = p_7 = 1.3p_{\circ,1}$, where p_i , for i = 1, ..., 8, is the pressure in cell *i*, and f = 0

MP, B. Seguin, Bull Math Biology, 2016

Intercellular transport of signalling molecules

Signalling molecules interact with cells

- as a ligand for membrane receptors
- and/or by entering into the cell through its membrane or endocytosis

$$C+R_f \stackrel{b_e}{\rightleftharpoons}_{a_e} R_b$$

- Diffusion of signalling molecules c and s in the extra- and intracellular spaces
- Diffusion of free and bound receptors r_f and r_b and of active and inactive co-receptors (proteins) p_a and p_d on cell membrane
- Ligands c bind to r_f to produce r_b and s interact with p_a
- Bound receptors r_b dissociate into free receptors and ligands
 Active co-receptors (proteins) p_a
 dissociate into inactive co-receptors
 (proteins) and bound receptors

Mathematical model for intercellular signalling

Diffusion, production and decay of ligands in extracellular space

 $\begin{array}{ll} \partial_t c = \nabla \cdot (D_e(x) \nabla c) + F_e(c) & \text{ in } \Omega_e, \ t > 0 \\ D_e(x) \nabla c \cdot \nu = 0 & \text{ on } \partial \Omega, \ t > 0 \\ c(0) = c_0 & \text{ in } \Omega_e \end{array}$



t, t > 0

• Equations for the receptors on the cell surface Γ

$$\partial_t r_f = D_f \Delta_{\Gamma} r_f + F_r(r_f, r_b) - a_e(x) r_f c + b_e(x) r_b - d_f r_f \quad \text{on } \Gamma$$
$$\partial_t r_b = D_b \Delta_{\Gamma} r_b \qquad + a_e(x) r_f c - b_e(x) r_b - d_b r_b \quad \text{on } \Gamma$$

Binding on the cell surfaces

$$D_e(x) \nabla c \cdot \nu = -a_e(x) c r_f + b_e(x) r_b \quad \text{on } \Gamma$$

Activation of an intracellular signalling pathway by r_b

 c, r_f, r_b density of ligands/receptors, d_f, d_b rate of decay of ligands/receptors F_e, F_r product. of ligands/receptors, D_e, D_f, D_b diffusion coefficients

Microscopic model for signalling processes

Diffusion, production and decay of signalling molecules

$$\partial_t c = \nabla \cdot (D_e^{\varepsilon}(x)\nabla c) + F_e(c) \quad \text{in } \Omega_e^{\varepsilon}, \ t > 0$$

 $\partial_t s = \varepsilon^2 \nabla \cdot (D_i^{\varepsilon}(x)\nabla s) + F_i(s) \quad \text{in } \Omega_i^{\varepsilon}, \ t > 0$

• Equations for the receptors /proteins on the cell surface Γ^{ε}

$$\partial_t r_f = \varepsilon^2 D_f \Delta_{\Gamma} r_f - G_e(c, r_f, r_b) + F_r(r_f, r_b) - d_f r_f$$

$$\partial_t r_b = \varepsilon^2 D_b \Delta_{\Gamma} r_b + G_e(c, r_f, r_b) - G_d(r_b, p_d, p_a) - d_b r_b$$

$$\partial_t p_d = \varepsilon^2 D_d \Delta_{\Gamma} p_d - G_d(r_b, p_d, p_a) + F_d(p_d) - d_d p_d$$

$$\partial_t p_a = \varepsilon^2 D_a \Delta_{\Gamma} p_a + G_d(r_b, p_d, p_a) - G_i(p_a, s) - d_a p_a$$

 \blacktriangleright Binding on the cell surfaces Γ^{ε}

$$D_e^{\varepsilon}(x) \nabla c \cdot \nu = -\varepsilon G_e(c, r_f, r_b) \quad \text{on } \Gamma^{\varepsilon}, \ t > 0$$

$$\varepsilon^2 D_i^{\varepsilon}(x) \nabla s \cdot \nu = \varepsilon G_i(p_a, s) \quad \text{on } \Gamma^{\varepsilon}, \ t > 0$$

 $\overbrace{Y_e}^{\Gamma}$

Binding reactions

$$G_{e}(c, r_{f}, r_{b}) = a_{e}^{\varepsilon}(x) c r_{f} - b^{\varepsilon}(x) r_{t}$$
$$G_{d}(r_{b}, p_{d}, p_{a}) = a_{i}(x) r_{b} p_{d}$$
$$G_{i}(p_{a}, s) = \gamma_{i}^{\varepsilon}(x) p_{a} - \bigvee_{Y_{e}}^{\Gamma}$$

Macroscopic equations

Macroscopic concentrations

$$\partial_t c - \nabla \cdot (D_e^{\text{hom}} \nabla c) = F_e(c) - \frac{1}{|Y_e|} \int_{\Gamma} G_e(c, r_f, r_b) \, d\gamma_y \quad \text{in } \Omega_T$$
$$\partial_t s - \nabla_y \cdot (D_i(y) \nabla_y s) = F_i(s) \quad \text{in } \Omega_T \times Y_i$$

► Receptors distribution on the cell surface on $\Omega_T \times \Gamma$

$$\partial_t r_f = D_f \Delta_{\Gamma,y} r_f - G_e(c, r_f, r_b) + F_r(r_f, r_b) - d_f r_f$$

$$\partial_t r_b = D_b \Delta_{\Gamma,y} r_b + G_e(c, r_f, r_b) - G_d(r_b, p_d, p_a) - d_b r_b$$

+ equations for p_a, p_d

Macroscopic coefficients

$$D_{e,ij}^{\text{hom}}(x) = \frac{1}{|Y_e|} \sum_{k=1}^{3} \int_{Y_e} (D_{e,ij}(x,y) + D_{e,ik}(x,y) \partial_{y_k} w_j) \, dy$$

where

$$\begin{aligned} -\nabla_y \cdot (D_e(x, y)(\nabla_y w^j + e_j)) &= 0 \text{ in } Y_e, \\ -D_e(x, y)(\nabla_y w^j + e_j) \cdot \nu &= 0 \text{ on } \Gamma, \quad w^j \quad Y - \text{periodic} \end{aligned}$$

 $\Omega_{\mathcal{T}} = (0, \mathcal{T}) imes \Omega$



Multiscale numerical simulations

$$\operatorname{div}_{y}(D_{e}^{*}(\nabla_{y}w^{j}+e_{j})) = 0 \quad \text{in } Y_{e}, \quad \int_{Y_{e}} w^{j}(y)dy = 0,$$
$$D_{e}^{*}(\nabla_{y}w^{j}+e_{j}) \cdot \nu = 0 \quad \text{on } \Gamma, \quad w^{j} \quad Y - \text{periodic},$$



$$D_{h,e}^{\text{hom}} = \begin{bmatrix} 8.167 \cdot 10^{-3} & 0\\ 0 & 1.841 \cdot 10^{-3} \end{bmatrix} \qquad D_{h,e}^{\text{hom}} = \begin{bmatrix} 6.556 \cdot 10^{-3} & 0\\ 0 & 6.149 \cdot 10^{-3} \end{bmatrix}$$

$$D_e^* = 10^{-2}, \quad D_i^* = 10, \quad D_f^* = D_b^* = D_d^* = D_a^* = 10^{-2},$$

MP, C. Venkataraman, arXiv 2018

Effect of anisotropic microstructure

